

# 极化算符

$$\vec{P} = \int d^3r \vec{r} \hat{P}(\vec{r})$$

$$\frac{\partial \vec{P}}{\partial t} = \int d^3r \vec{r} \frac{\partial \hat{P}(\vec{r})}{\partial t}$$

由于  $\frac{\partial \hat{P}(\vec{r})}{\partial t} = -\nabla \cdot \vec{J}$

因此

$$\frac{\partial \vec{P}}{\partial t} = -\int d^3r \vec{r} \nabla \cdot \vec{J}$$

由于  $\nabla \cdot (\vec{J} \vec{r}) = (\nabla \cdot \vec{J}) \vec{r} + \vec{J} \cdot \nabla \vec{r}$

因此  $\frac{\partial \vec{P}}{\partial t} = -\int d^3r \nabla \cdot (\vec{r} \vec{J}) + \int d^3r \vec{J} \cdot \nabla \vec{r}$

由于无穷远处电流为0

则  $\int d^3r \nabla \cdot (\vec{r} \vec{J}) = 0$

$$\frac{\partial \vec{P}}{\partial t} = \int d^3r \vec{J} \cdot \nabla \vec{r} = \int d^3r \vec{J}$$

# 极化算符离散化

由于  $\hat{P}(\vec{r}) = \sum_i \delta(\vec{r} - \vec{r}_i)$ , 二次量子化后

$\hat{P}(\vec{r}) = \psi^\dagger(\vec{r}) \psi(\vec{r})$  是粒子数密度算符

离散化  $\hat{P}(\vec{r}) \Rightarrow \hat{n}_i = a_i^\dagger a_i$

$$\vec{r} \Rightarrow \vec{R}_i$$

$$\vec{p} = \sum_i \vec{R}_i \hat{n}_i$$

$$\text{由于 } \frac{\partial \vec{p}}{\partial t} = i [\hat{H}, \vec{p}]$$

$$\hat{H} = \underbrace{\sum_j \epsilon_j C_j^\dagger C_j}_{\text{在位项}} + \underbrace{\sum_{j\delta} W C_{j+\delta}^\dagger C_j}_{\text{跃迁项}} + \frac{1}{2} \underbrace{\sum_{ij} \hat{n}_i \hat{n}_j V_{ij}}_{\text{相互作用项}}$$

$$\hat{n}_i = C_i^\dagger C_i$$

$$[\hat{H}, \vec{p}] = \textcircled{1} \sum_j \sum_i \epsilon_j \vec{R}_i [C_j^\dagger C_j, C_i^\dagger C_i] + \textcircled{2} \sum_{j\delta} \sum_i W \vec{R}_i [C_{j+\delta}^\dagger C_j, C_i^\dagger C_i] \\ + \textcircled{3} \sum_{ij} \frac{1}{2} V_{ij} [\hat{n}_i \hat{n}_j, \hat{n}_i]$$

注意到，费米子满足反对易关系

$$[C_i, C_j^\dagger]_+ = \delta_{ij}, [C_i, C_j]_+ = 0, [C_i^\dagger, C_j^\dagger]_+ = 0$$

$$[C_j^\dagger C_j, C_i^\dagger C_i] = C_j^\dagger [C_j, C_i^\dagger C_i] + [C_j^\dagger, C_i^\dagger C_i] C_j \\ = C_j^\dagger \underbrace{[C_j, C_i^\dagger]_+}_{\delta_{ji}} C_i - C_i^\dagger \underbrace{[C_j^\dagger, C_i]_+}_{\delta_{ji}} C_j \\ = C_i^\dagger C_i - C_i^\dagger C_i \\ = 0$$

因此，第①项为0

$$\text{由于 } [C_j^\dagger C_j, C_i^\dagger C_i] = [\hat{n}_j, \hat{n}_i]$$

所以第③项也为0

$$\begin{aligned} \text{第②项: } [C_{j+\delta}^\dagger C_j, C_i^\dagger C_i] &= C_{j+\delta}^\dagger [C_j, C_i^\dagger C_i] + [C_{j+\delta}^\dagger, C_i^\dagger C_i] C_j \\ &= C_{j+\delta}^\dagger C_i \delta_{ij} - C_i^\dagger C_j \delta_{i, j+\delta} \end{aligned}$$

$$\begin{aligned} \text{因此 } [\hat{H}, \hat{P}] &= \sum_{j\delta} W (\vec{R}_j - \vec{R}_{j+\delta}) C_{j+\delta}^\dagger C_j \\ &= -\sum_{j\delta} W \vec{\delta} C_{j+\delta}^\dagger C_j \end{aligned}$$

$$\frac{\partial \vec{P}}{\partial t} = i[\hat{H}, \hat{P}] = -i \sum_{j\delta} W \vec{\delta} C_{j+\delta}^\dagger C_j$$

$$\text{所以 } \vec{J} = -iW \sum_{j\delta} \vec{\delta} C_{j+\delta}^\dagger C_j$$